

## Simulation of multi-step $\gamma$ cascade spectra from $^{59}\text{Fe}(d,p\gamma)$ reaction

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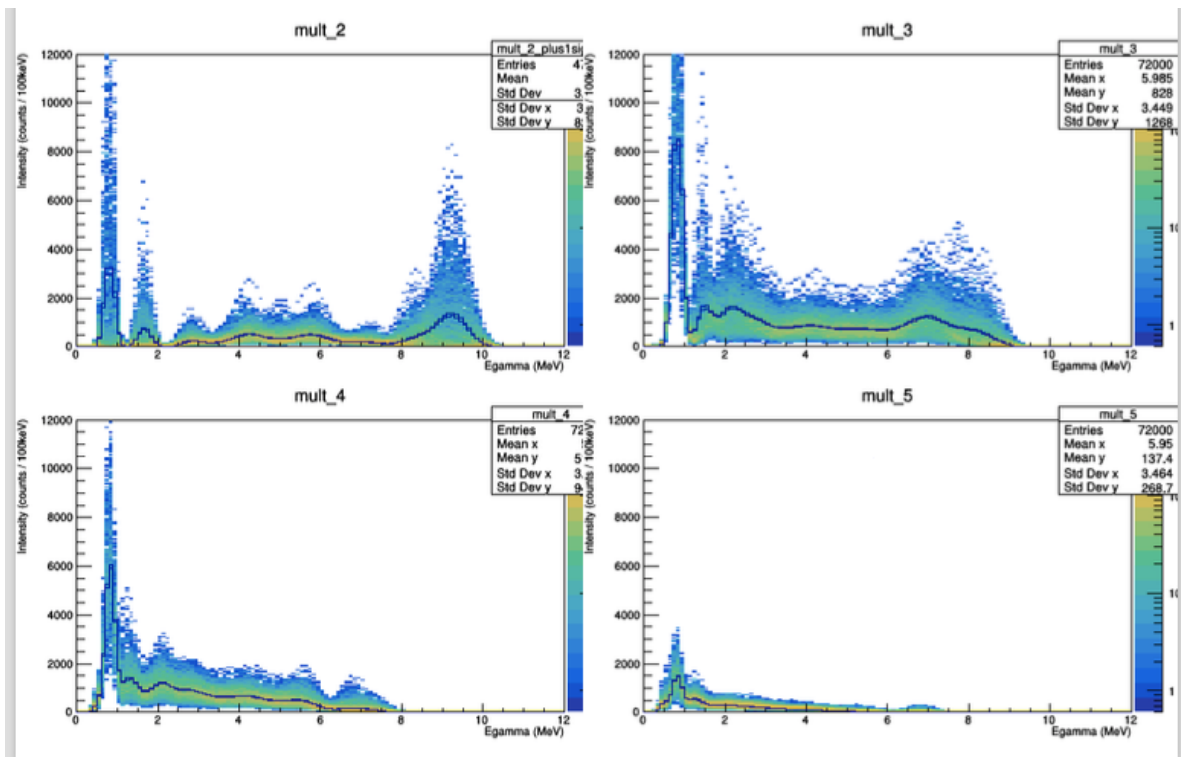
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Cosmic  $\gamma$  rays from long-lived ( $\tau$  (1/2)  $\sim$  Myrs) unstable nuclei such as  $^{26}\text{Al}$  and  $^{60}\text{Fe}$  are observed overall in the Galactic disk [1]. Understanding the origin of these nuclei will lead to greatly constrain and improve models of galactic chemical evolution and stellar nucleosynthesis.

These  $^{60}\text{Fe}$  are believed to be produced by the  $^{59}\text{Fe}(n,\gamma)$  reaction, e.g., in the He burning stage of massive stars. It is therefore important to determine the reaction cross sections at the stellar temperature (neutron energy ( $E_n$ )  $\sim$ 30 keV). However, since  $^{59}\text{Fe}$  is short-lived ( $\tau$  (1/2) = 44 days), indirect methods must be taken to determine the cross sections. We are developing a method using the  $^{59}\text{Fe}(d,p\gamma)^{60}\text{Fe}$  reaction to indirectly constrain the  $^{59}\text{Fe}(n,\gamma)$  cross sections. A charged particle detector with high energy-resolution (for ejectile protons) and a large  $\gamma$ -detector array with high detection efficiency are demanded for the method. Thus, a detector system composed of a Si detector and 128 BaF<sub>2</sub> crystals [2] is under development at the TAMU Cyclotron Institute [3].

Data analysis from the present method is highly similar to the one employed for (n, $\gamma$ ) experiments with the DANCE BaF<sub>2</sub> array at Los Alamos National Laboratory (LANL) [4]. In our case, highly excited states of a nucleus of interest are populated by ( $d, p$ ) reactions instead of neutron capture reactions. Populated states and underlying states of a nucleus,  $^{60}\text{Fe}$  in our case, are simulated by a Monte Carlo Simulation code named DICEBOX [5]. The simulation process, “Nuclear Realization”, involves multiple nuclear realizations to accommodate the fact that many energy levels and decay properties of  $^{60}\text{Fe}$  are not known in details. Realistic energy levels of  $^{60}\text{Fe}$  from the ground state to a certain excitation energy (usually set to the neutron separation energy, 8.8 MeV for  $^{60}\text{Fe}$ ), with e.g., given spins, parities,  $\gamma$  partial width, are thus generated based on a nuclear level density model and photon strength function models provided in the input file. The multi-step (MS)  $\gamma$  cascades from the excitation energy to ground state are incorporated in our GEANT4 simulation in which the TAMU BaF<sub>2</sub> array is built. The obtained GEANT4 spectra will be eventually compared with our experimental ( $d,p\gamma$ ) spectrum. By iterating the simulation varying the models used, we will find the “best-match” level density and PSF models with the experimental spectra.  $^{59}\text{Fe}(n,\gamma)$  cross sections can be determined using these models, e.g., with the TALYS code [6].

Figs. 1 show examples of simulated (generated) MS  $\gamma$  cascade spectra for  $^{58}\text{Fe}$ , which were made to study data from the test experiments with the  $^{57}\text{Fe}(d,p\gamma)$  reaction to benchmark the method. Note no detector responses have been included in the present simulations yet. The simulations are made from a nuclear level density available from [7] and PSFs theoretically calculated by Goriely *et al.* [8]. Each panel presents the overlaid spectra from different nuclear realizations for the cases of  $\gamma$  decay multiplicity



**Fig. 1.** Multi step  $\gamma$  cascade spectra from  $E_x=10$  MeV  $^{58}\text{Fe}$   $J=2^+$  state. Each panel shows a different  $\gamma$  source multiplicity.

equal to 2, 3, 4, and 5, respectively (see e.g., [9]). The lines in the panels correspond to average values at respective  $\gamma$  energies. It is clear the MS  $\gamma$  cascade spectra have intrinsic fluctuations derived from quantum mechanics (more specifically, e.g., the Porter-Thomas distribution). Therefore, a band corresponding to 66% (1 sigma) of the data will be extracted from each spectrum and then compared with the experimental data.

Fig. 2 shows the  $^{57}\text{Fe}(n,\gamma)$  cross sections calculated using the models above with TALYS in comparison to past measurements [10, 11]. The calculations are preliminary, given no appropriateness of models used have been discussed yet. It can be seen, however, the  $(n,\gamma)$  cross sections can vary by a factor of two easily depending on models used. Thus, it is essential for the TAMU  $\text{BaF}_2$  array to have the capability to distinguish different nuclear level density models and PSF models from the measured spectra. It is also noteworthy that the experimental cross-sections in the energy range of astrophysical interest (the range shown in the Fig. 2) have some discrepancies from different measurements. PSFs of  $^{58}\text{Fe}$  currently being studied at LANL [12] may instead be a better guide for us to benchmark our method instead.

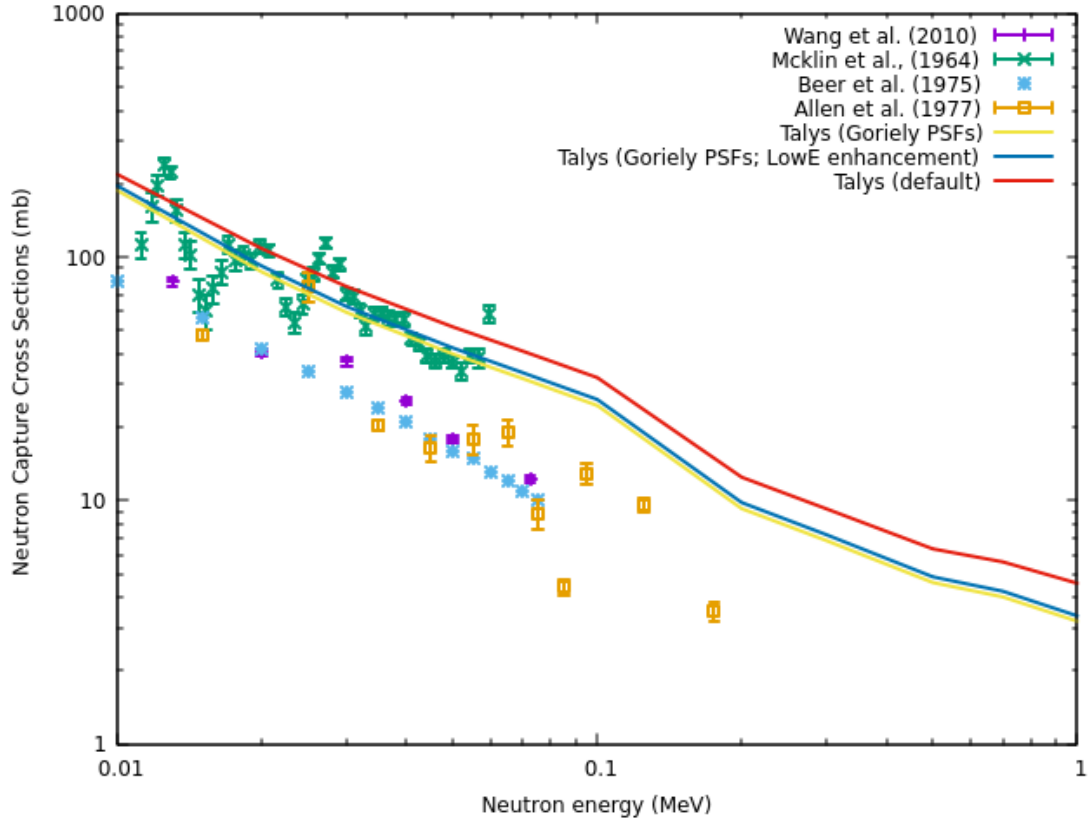


Fig. 2.  $^{57}\text{Fe}(n,\gamma)$  cross sections from the past measurements and preliminary TALYS calculations.

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